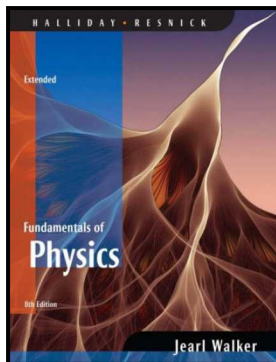


Workshop Physics

1017 - 312

# University Physics II



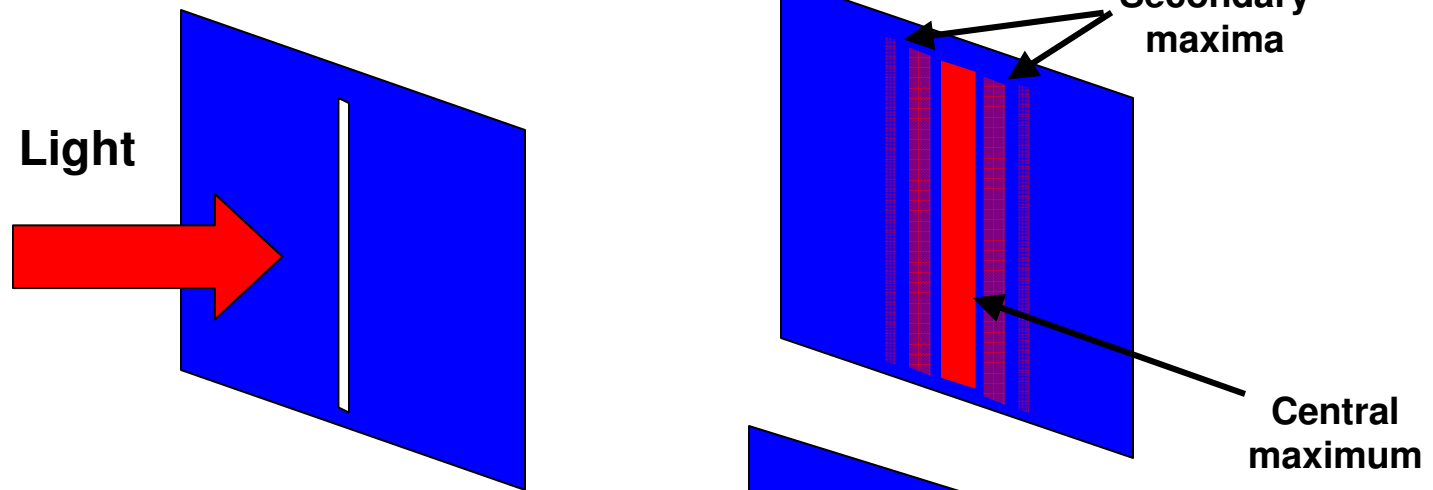
**Week 9 : Day 2**

# Outline

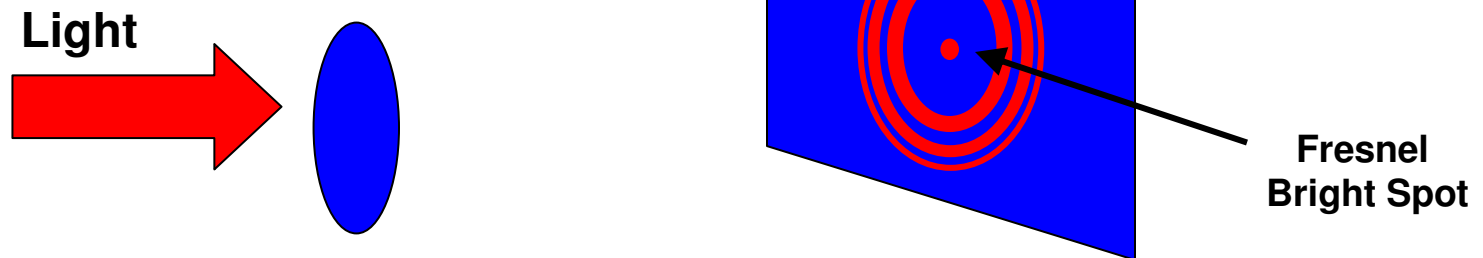
- ❑ **The wave theory of light**
  - Diffraction
  - Interference
- ❑ **Single-slit diffraction**
  - Locating diffraction minima
  - Intensity pattern
- ❑ **Activity - Diffraction**

# Diffraction and the Wave Theory of Light

Diffraction from a single narrow slit:

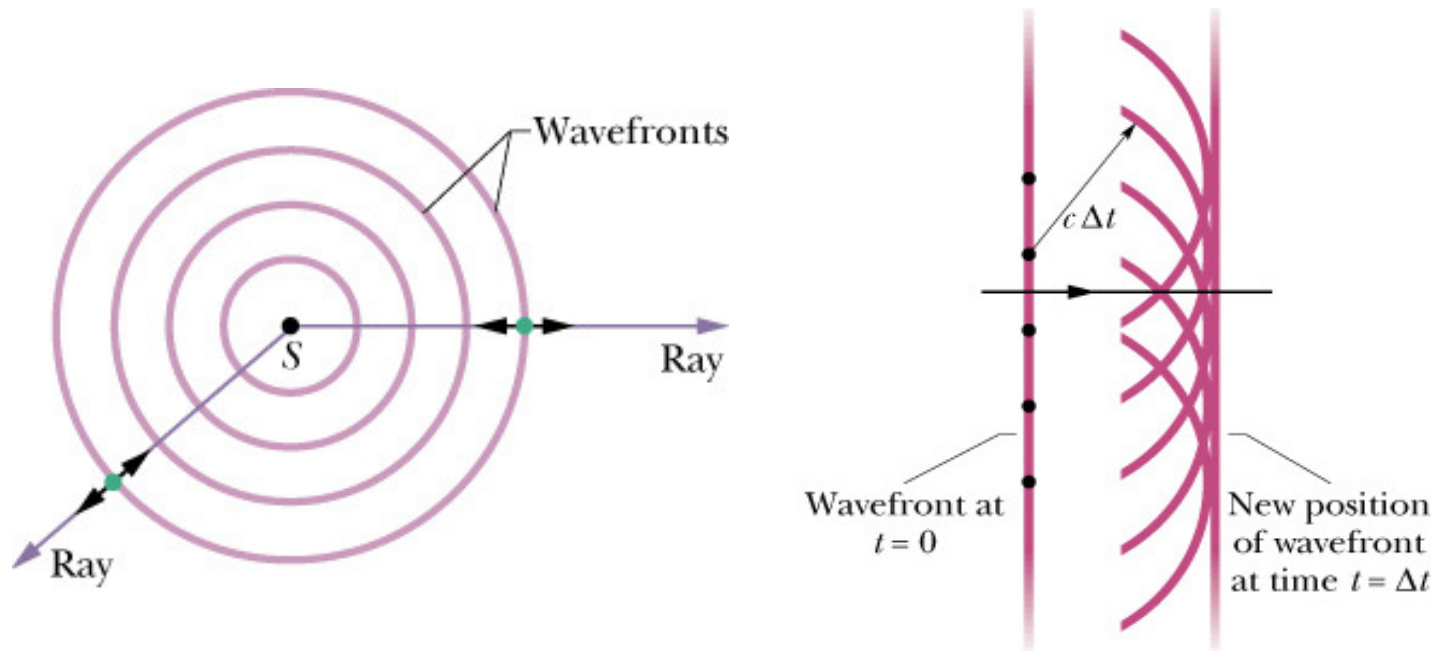


Diffraction from a lens:



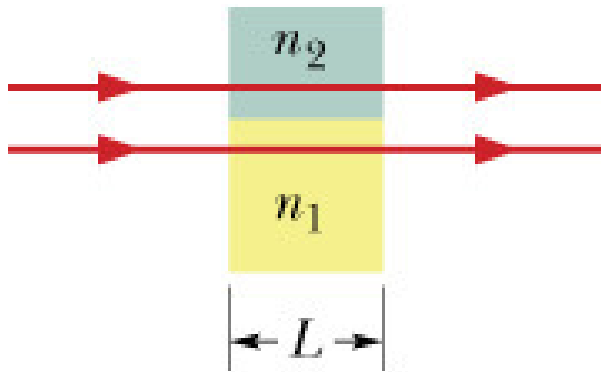
# Light as a Wave

Huygen's Principle: All points on a wavefront serve as point sources of spherical secondary wavelets.



After time  $t$ , the new position of the wavefront will be that of a surface tangent to these secondary wavelets.

## Wavelength and the Index of Refraction



The frequency of light in a medium is the same as it is in vacuum.

$$\lambda_n = \frac{\lambda}{n} \quad f_n = \frac{v}{\lambda_n} = \frac{c/n}{\lambda/n} = \frac{c}{\lambda} = f$$

Since wavelengths in  $n_1$  and  $n_2$  are different, the two beams may no longer be in phase.

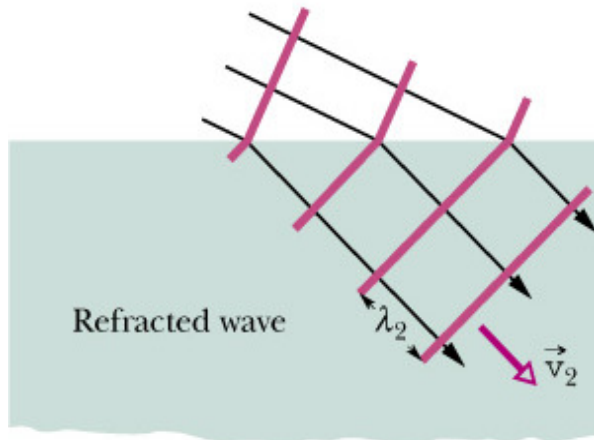
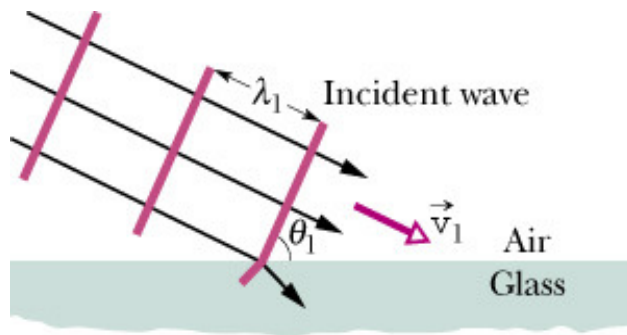
Number of wavelengths in  $n_1$ :  $N_1 = \frac{L}{\lambda_{n1}} = \frac{L}{\lambda/n_1} = \frac{Ln_1}{\lambda}$

Number of wavelengths in  $n_2$ :  $N_2 = \frac{L}{\lambda_{n2}} = \frac{L}{\lambda/n_2} = \frac{Ln_2}{\lambda}$

Assuming  $n_2 > n_1$ :  $N_2 - N_1 = \frac{Ln_2}{\lambda} - \frac{Ln_1}{\lambda} = \frac{L}{\lambda}(n_2 - n_1)$

⇒  $N_2 - N_1 = 1/2$  wavelength → destructive interference

# Wavelength and the Law of Refraction

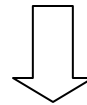


Law of Refraction:

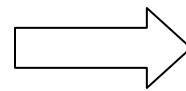
$$n_1 \sin \theta_1 = n_2 \sin \theta_2$$

Index of Refraction:

$$n_1 = \frac{c}{v_1} \quad \text{and} \quad n_2 = \frac{c}{v_2}$$



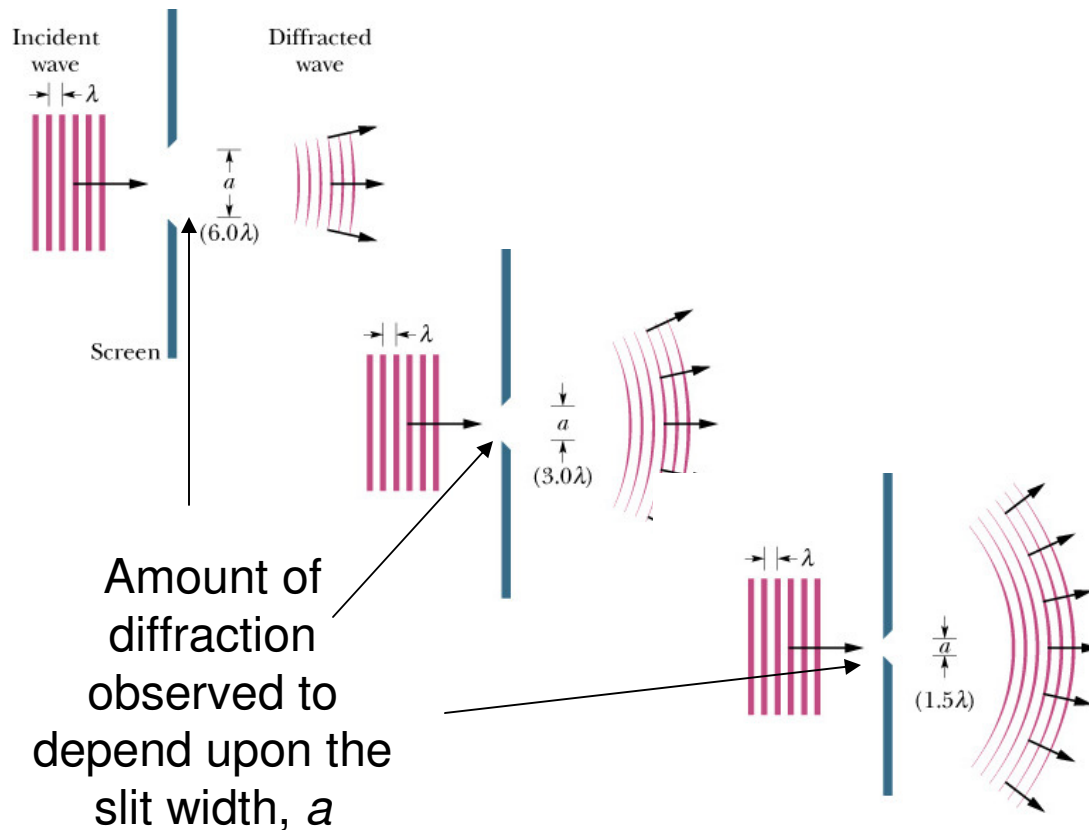
$$n = \frac{c}{v} = \frac{c}{f\lambda_n} = \frac{c}{f} \frac{1}{\lambda_n}$$



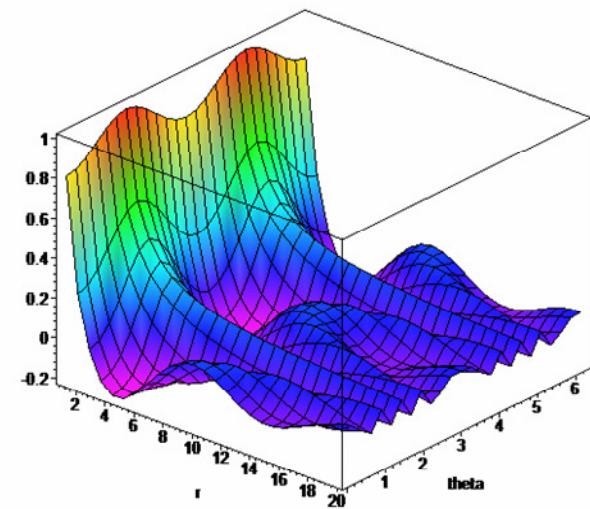
$$\frac{\sin \theta_1}{\sin \theta_2} = \frac{\lambda_1}{\lambda_2}$$

# Single-Slit Diffraction

For plane waves entering a single slit, the waves emerging from the slit start spreading out or *diffracting*:

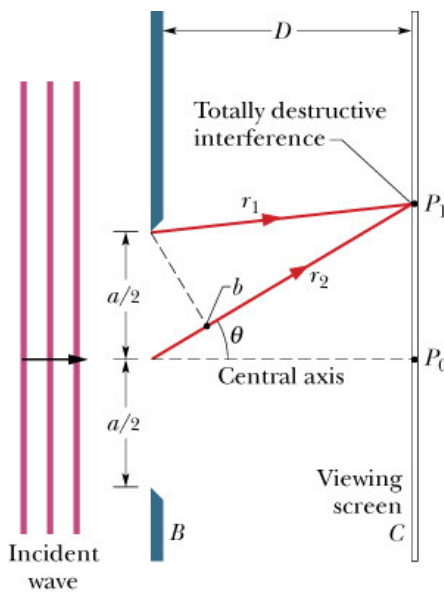


A Plane Wave

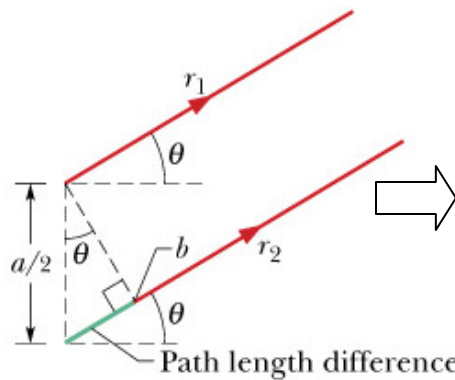


$$\psi(r, \theta) = A \frac{e^{-2\pi i k r \cos \theta}}{r}$$

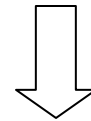
## Locating Single-Slit Diffraction Minima



Setting path length difference to  $\lambda/2$  for each pair of rays, we obtain the first dark fringes at:



$$\left( \begin{array}{c} \text{phase} \\ \text{difference} \end{array} \right) = \left( \frac{2\pi}{\lambda} \right) \left( \begin{array}{c} \text{path length} \\ \text{difference} \end{array} \right)$$

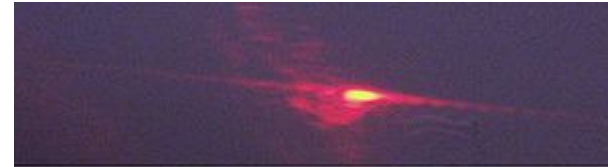


(first minimum)  $\frac{a}{2} \sin \theta = \frac{\lambda}{2} \rightarrow a \sin \theta = \lambda$

Dividing the slit into increasingly larger even numbers of zones, we can find higher order minima:

$$a \sin \theta = m\lambda, \text{ for } m = 1, 2, 3 \dots$$

# Single-Slit Diffraction Intensity



Pattern produced from a single slit.

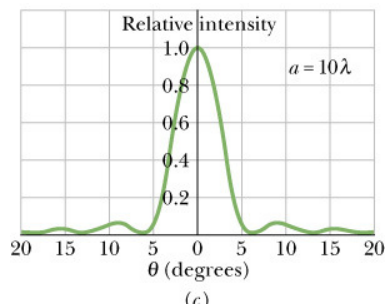
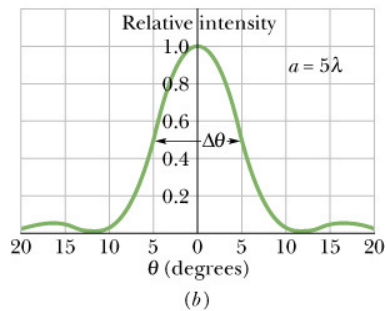
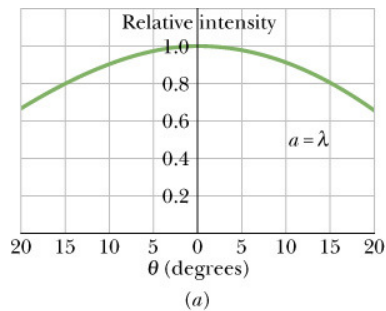
The intensity at the screen due to a single slit is:

$$I(\theta) = I_m \left( \frac{\sin \alpha}{\alpha} \right)^2$$

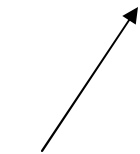
where  $\alpha = \frac{1}{2} \phi = \frac{\pi a}{\lambda} \sin \theta$

Minima in the intensity occur when:

$$a \sin \theta = m\lambda, \quad \text{for } m = 1, 2, 3 \dots$$



Width of intensity inversely proportional to slit width,  $a$

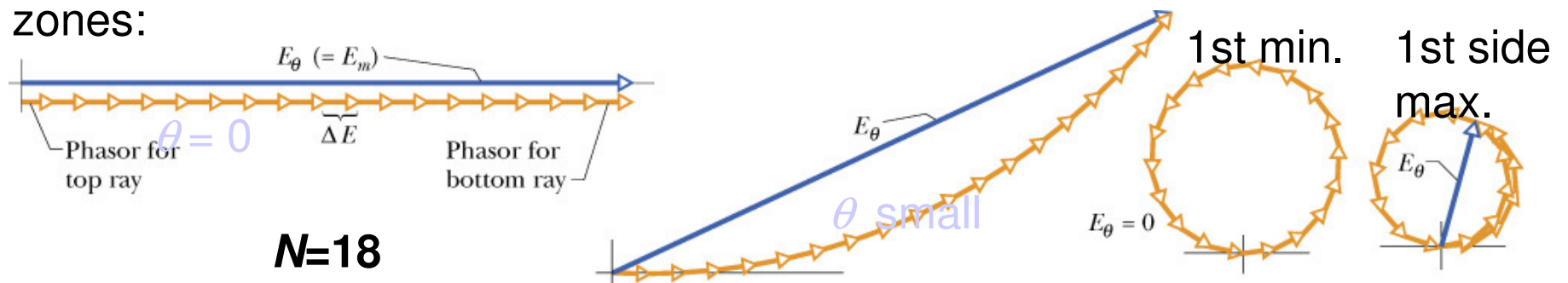


## Intensity in Single-Slit Diffraction, Qualitatively

To obtain the locations of the minima, the slit was equally divided into  $N$  zones, each with width  $\Delta x$ . Each zone acts as a source of Huygens wavelets. Now these zones can be superimposed at the screen to obtain the intensity as a function of  $\theta$ , the angle to the central axis.

$$\left( \begin{array}{c} \text{phase} \\ \text{difference} \end{array} \right) = \left( \frac{2\pi}{\lambda} \right) \left( \begin{array}{c} \text{path length} \\ \text{difference} \end{array} \right) \implies \Delta\phi = \left( \frac{2\pi}{\lambda} \right) (\Delta x \sin \theta)$$

To find the net electric field  $E_\theta$  (intensity  $\propto E_\theta^2$ ) at point  $P$  on the screen, we need the phase relationships among the wavelets arriving from different zones:



# Activity – Diffraction

## □ Examine single-slit parameters

- Width of central bright fringe
- Width of secondary bright fringe

Your Name (Print): \_\_\_\_\_ Date: \_\_\_\_\_  
 Group Members: \_\_\_\_\_ Group: \_\_\_\_\_

### Diffraction

**Caution:** Class IV lasers (mostly harmless) are used in this experiment. Under no circumstances should you look directly into the laser, even if it seems like a really good idea at the time and your lab partners bet you a quarter it won't hurt. Laser light reflected from non-metallic surfaces is safe, but be careful with reflective surfaces. Please help save batteries and turn them off if you're doing calculations for any length of time.

The single-slit diffraction geometry is shown in Figure 1. As discussed in class, the first position of fully destructive interference (the center of a dark fringe) occurs when the light from the top half of a slit of width  $a$  interferes destructively with the light from the bottom half of the slit. In general, the angular positions of the dark fringes  $\theta_m$  are determined by

$$a \sin \theta_m = m\lambda, \text{ where } m = \pm 1, \pm 2, \pm 3, \dots \text{ but not } m = 0 \quad (1)$$

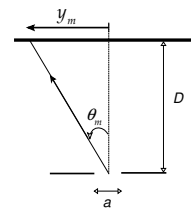
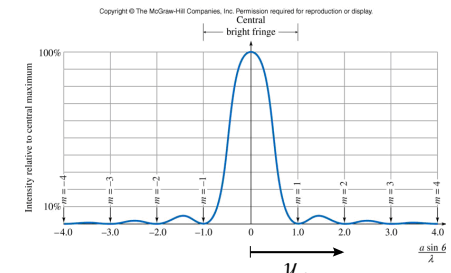


Figure 1 Single-slit Diffraction Geometry



GRR 2nd Figure 25.31 b

Instead of measuring the angle  $\theta_m$  directly, you will measure the linear distance  $y_m$  along the screen from the center of the pattern ( $\theta = 0$ ) to the center of the  $m^{\text{th}}$  dark fringe. The angle  $\theta_m$  is related to the distance  $y_m$  and the distance  $D$  from the slit to the screen; specifically,

$$\tan \theta_m = \frac{y_m}{D} \quad (2)$$

If  $\theta_m$  is small, then

$$\sin \theta_m \approx \tan \theta_m = \frac{y_m}{D} \quad (3)$$